

# Event shape engineering via Glauber MC model

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**Abstract.** In our exploration of event shape engineering, the Glauber model served as a foundational tool for understanding the anisotropic geometry of the Quark-Gluon Plasma (QGP). Utilizing the TGlauberMC-3.2\* model within ROOT, we systematically analyzed one million events. From the  $\epsilon_2$  &  $N_{part}$  plot, our data revealed an average maximum  $dN$  value of 12.00 with associated parameters:  $\epsilon_2 = 0.91$ ,  $\psi_2 = 2.74$ ,  $\psi_3 = 0.91$  and  $N_{part} = 14.00$ . These findings illuminate the distinct configurations that yield the most pronounced anisotropic geometries of the QGP, providing insights into optimizing event shape configurations.

**Keywords:** Event Shape Engineering, Glauber Model, Quark-Gluon Plasma (QGP), TGlauberMC-3.2, Anisotropic Geometries.

## 1. Introduction

### 1.1. QGP and event shape engineering

Quark gluon plasma is the state of matter, consisting of free gluons and quarks, produced in ultra relativistic heavy ion collisions with extremely high temperature and density. It's believed that QGP existed shortly after the big bang, thus by studying the QGP, valuable insights can be gained to better understand the fundamental properties of matter and early formation of the universe.

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\* (<http://www.hepforge.org/downloads/tglaubermc>) for the most recent TGlauberMC release (currently version 3.2).

### 1.2. The Glauber Monte Carlo model and path length measurement

Event shape engineering is the technique to manipulate the shapes of QGP and extract the property of produced particles. In this letter we focus on the anisotropy flow that provide insight into the initial and subsequent evolution of the collision system. By disentangling various flow components, the dynamics of the created QGP can be shown.

In heavy ion collisions, QGP is formed in the overlap region of the colliding nucleus. We use the Glauber model to describe the collision. It is assumed that the nucleus-nucleus collision can be simplified to uncorrelated nucleon-nucleon collisions.[1]

The position of each nucleon is sampled from the Fermi distribution [2]. Collisions can be simulated with either a random or fixed impact parameter and then projected onto the x-y plane. In the Monte Carlo approach, a nucleon-nucleon collision occurs if the distance between the nucleons in the x-y plane satisfies the condition given by

$$r = \sqrt{\frac{\text{cross section}}{\pi}}.$$

The TGlauberMC code, integrated within Root, is employed to execute the Monte Carlo simulations and simulate collisions. The code produces data related to collision positions and other pertinent parameters, such as the impact parameter  $b$ , elliptical eccentricity  $\epsilon_2$ , triangular eccentricity  $\epsilon_3$  [3], number of collisions  $N_{\text{coll}}$ , and number of participants  $N_{\text{part}}$ .

The study of the QGP often evolves the anisotropy of its shape. In this experiment we use the path length difference to manifest the anisotropy. Paths are possible route traversed by partons. [4] Due to the high density of the QGP, the path length is measured by counting the number of nucleons along the path. The path length difference is defined as the length difference between minor axis and major axis. The direction of the minor axis is determined by angle  $\psi_2$  and the direction of major axis is perpendicular to the minor axis.

## 2. Methodology and approach

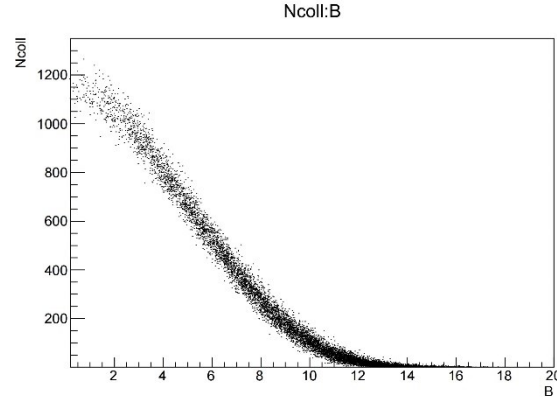
### 2.1. C++ based ROOT frame

ROOT frame is a framework used to process data first born at CERN. The frame work is dominant to the researches of high energy physics.[5] In this paper, C++ based ROOT frame is used to run the Monte Carlo Glauber (MCG) model of version 3.2 to simulate  $Au + Au$  collisions and generate data of these collisions, the model is run directly on terminal of MacOS in this project.

### 2.2. Visualisation of data

#### 2.2.1. Verification of generated data

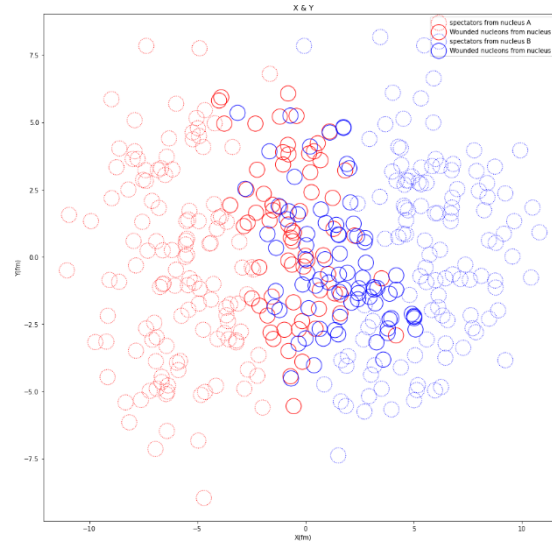
Function `runAndSaveNtuple()` generates a series of different MCG events and saves related physical quantities.[1] In this paper, cross section is set to 40mb, which corresponds to collision energy of 200Gev and nucleon radius of 0.564fm. The graph relating  $N_{\text{coll}}$  and impact parameter is plotted using ROOT shown in Figure 1. The shape and values of the graph match with the result in Reference [2], which verifies the process of generating data using the MCG model.



**Figure 1.** Illustration of the relationship between the number of collisions  $N_{\text{part}}$  and the impact parameter  $B$  over 10,000 elastic  $Au + Au$  collisions.

### 2.2.2. Visualisation of the Quark-Gluon Plasma (QGP)

The coordinates and number of collisions on an individual nucleon during one  $Au + Au$  event with cross section  $40mb$  is calculated by the MCG model. The typical shapes of collisions can be seen By plotting the positions of nucleons in the collisions. In Figure 2, a collision of impact parameter 8 is shown, with red and blue dotted-line circles representing the spectators of nucleus A and nucleus B, and red and blue full-line circles representing the wounded nucleons (nucleons that participate in the collision) of nucleus A and nucleus B.



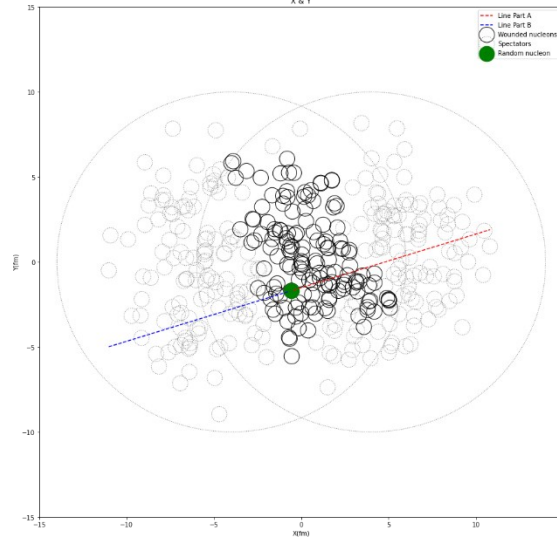
**Figure 2.** Representation of the Quark-Gluon Plasma in the  $XY$  plane.

### 2.2.3. Correlation between $N_1$ and $N_2$

In this step, an arbitrary wounded nucleon is chosen. A line with random direction in the  $XY$  plane passing through the random nucleon is divided by it into two segment:  $L_1$  and  $L_2$ . Wounded nucleons whose distance to the line is smaller than the radius of  $Au$  nucleons ( $r=0.564fm$ ) can be considered as passing through the line. The number of wounded nucleons passing through  $L_1$  is defined as  $N_1$ , and the number of wounded nucleons passing through  $L_2$  is defined as  $N_2$ .

In Figure 4, the plots of  $N_1$  and  $N_2$  of different impact parameters are illustrated. A strong anti-correlation between  $N_1$  and  $N_2$  can be seen clearly. The average value of both  $N_1$  and  $N_2$  are decreasing as  $B$  increases, which is as predicted due to a decrease of average number of participants with the increase of  $B$ . Possible errors can contribute to this result such as when no nucleon is passing through

the line with a random angle, and the possibility of fetching wounded nucleons on the edge of the QGP. To avoid the latter error, the process of fetching a random point is weighted using the  $N_{\text{coll}}$  of nucleons.



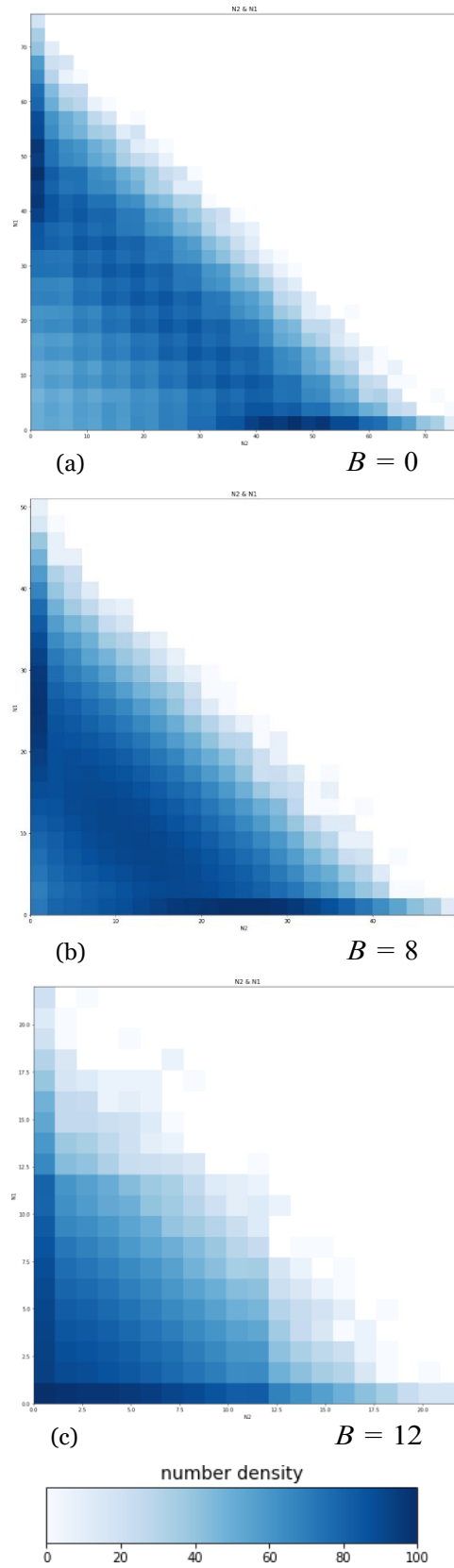
**Figure 3.** Computation of  $N_1$  and  $N_2$ . The segment  $L_1$ , depicted in red, represents points possessing x-coordinates greater than the designated random point. Conversely, the segment  $L_2$ , illustrated in blue, encompasses points with x-coordinates less than this random point. In the provided illustration, for instance, the values are determined as  $N_1 = 20$  and  $N_2 = 5$ .

#### 2.2.4. Event shape engineering for ellipse

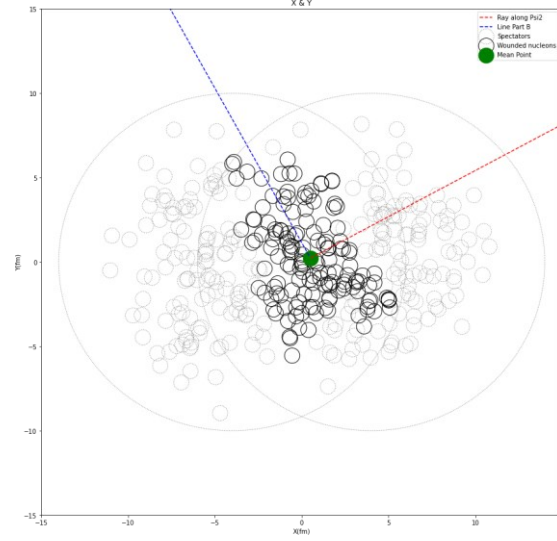
In our investigation, we introduce  $N_1$  and  $N_2$  as the quantities representing the number of nucleons along the major and minor axes, respectively. As illustrated in Figure 5,  $N_2$  corresponds to the count of wounded nucleons intersecting with the ray  $L_2$ . This ray, depicted in red within Figure 5, is conceptualized as a vector in a unit circle for the purposes of this algorithm. It is angled at  $\psi_2$ , which signifies the inclination of the minor axis of an ellipse with respect to the horizontal axis. For the specific event displayed in Figure 5,  $N_2$  is found to be 9.

On the other hand,  $N_1$  is associated with the number of wounded nucleons traversing the ray  $L_1$ . This ray, visualized in blue in Figure 5, is perpendicular to  $L_2$ . Importantly, it encompasses a greater number of nucleons than the ray situated in the opposing direction. For the event under consideration,  $N_1$  is ascertained to be 24. The Perpendicular Detector Algorithm was employed to determine the values of  $N_1$  and  $N_2$ .

For event shape engineering, our endeavor was centered around identifying an ellipse exhibiting optimal symmetry, and pinpointing an ellipse characterized



**Figure 4.** Depiction of the relationship between  $N_1$  and  $N_2$  for impact parameters of 0, 8, and 12.

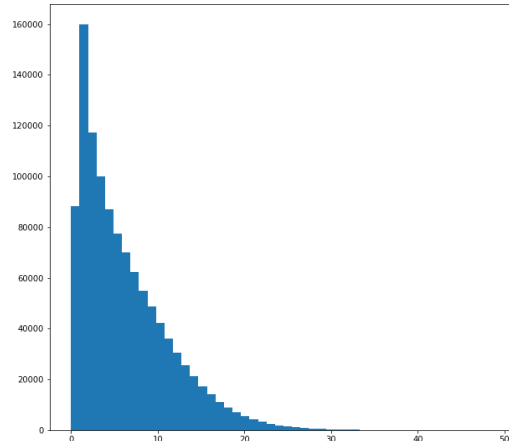


**Figure 5.** Depiction of the major and minor axes. The blue line signifies the major axis, while the red line denotes the minor axis oriented in the direction of  $\psi_2$ .

by the most marked difference in its dimensions, represented as  $N_1 - N_2$  or  $dN$ . After analyzing 1,102,826 events, we secured equivalent sets of  $N_1$  and  $N_2$  values. The distribution of  $dN$  is charted in Figure 6. For an in-depth scrutiny, we selected regions manifesting a pronounced  $dN$ , specifically in cases where  $dN > 15$ .

For data with  $dN > 15$ , contour lines are plotted on histograms. Figure 7(a) is for  $\epsilon_2$  (used to determine the level of symmetry of a shape to an ellipse) and  $N_{part}$ , Figure 7(b) is for  $\epsilon_3$  (used to determine the level of symmetry of a shape to a triangle) and  $N_{part}$ .

3-dimensional graphs are also plotted to show the correlation between those physical quantities and average values of  $dN$  by using surface 3d plotting tools from Matplotlib. For calculating the average values of  $dN$ , a histogram of a physical quantity chosen from  $\epsilon_2$ ,  $\epsilon_3$ ,  $\psi_2$ ,  $\psi_3$  as the y-axis and  $N_{part}$  as the x-axis is drawn with weight 1, and a histogram of the same quantities is drawn with weight  $dN$ . By dividing the two histograms, the average values of  $dN$  is represented as the colour depth of the histogram. Then the colour depth (average  $dN$ ) is used as the third axis (z-axis) of the surface 3d graphs. In Figure 9(a) and 9(b), surface 3d graphs are plotted for  $\epsilon_2$  and  $\epsilon_3$  being the y-axis respectively, and  $N_{part}$  for x-axis, average  $dN$  for z-axis.



**Figure 6.** Presentation of the  $dN$  distribution, with values extending from 0 to 30. For the Quark-Gluon Plasma (QGP), a substantial path difference is inferred when  $dN$  surpasses 15.

### 3. Result

In a detailed examination of the anisotropy of the Quark-Gluon Plasma (QGP), several intriguing observations can be made. Referring to Fig. 9(a), there is a notable increase in the path difference with  $\epsilon_2$ . This increase reaches its maximum, approximately at 12.00, when the number of participating nucleons,  $N_{\text{part}}$ , is 14.00. Conversely, as depicted in Fig. 9(b), the path difference remains substantial even for small values of  $\epsilon_3$ .

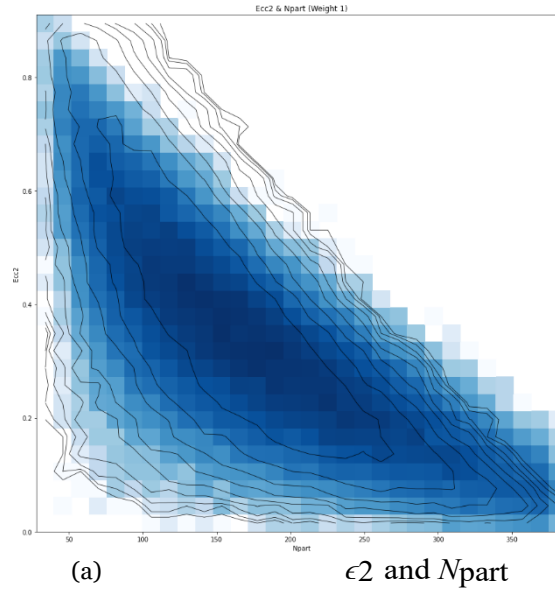
The optimal combination resulting in the largest path difference is achieved when  $\psi_2 = 2.74$ ,  $\psi_3 = 0.91$ , and  $\epsilon_2 = 0.91$ . This suggests that the anisotropy is at its zenith when the collision is peripheral and the number of participants is minimal. For a central collision,  $N_{\text{part}}$  approaches 400, corresponding to  $\epsilon_2 = 0.22$ .

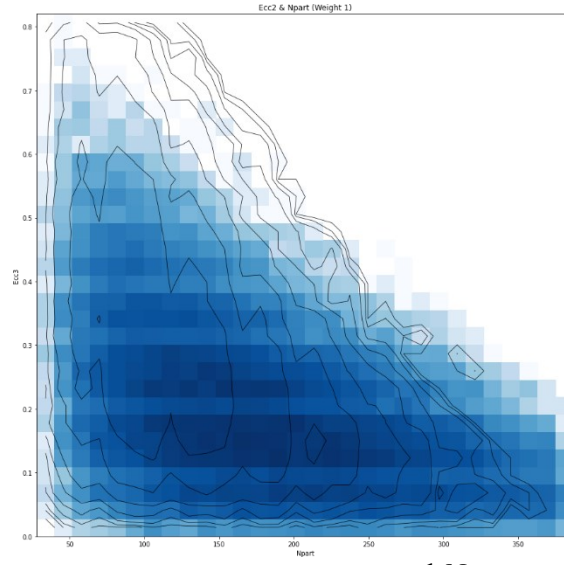
Interestingly, when  $N_{\text{part}}$  falls below 10, the QGP exhibits maximum asymmetry. However, due to the inherent limitation that  $dN$  cannot surpass  $N_{\text{part}}$ , this observation is bounded by the relatively small value of  $N_{\text{part}}$ .

### 4. Conclusion

In the present study, emphasis has been laid on modulating the path length difference of the Quark-Gluon Plasma (QGP) employing the Glauber model. The optimal set of parameters yielding the maximal path length difference comprises  $\epsilon_2 = 0.91$ ,  $\psi_2 = 2.74$ ,  $\psi_3 = 0.91$ , and  $N_{\text{part}} = 14.00$ . This particular combination is indicative of a highly peripheral collision. The magnitude of the path length difference is intrinsically contingent upon the size of the QGP, represented by the number of participating nucleons, as well as its anisotropic nature.

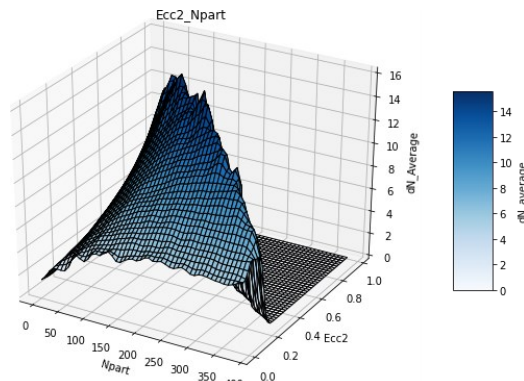
The pinnacle of path length difference can be realized only in the presence of an adequate number of participating nucleons in conjunction with a pronounced anisotropic geometry of the QGP. To bolster the veracity of the path length distribution, further simulation of data is recommended. Such an approach would not only augment the reliability but also ameliorate the effects of anomalies stemming from random fluctuations.



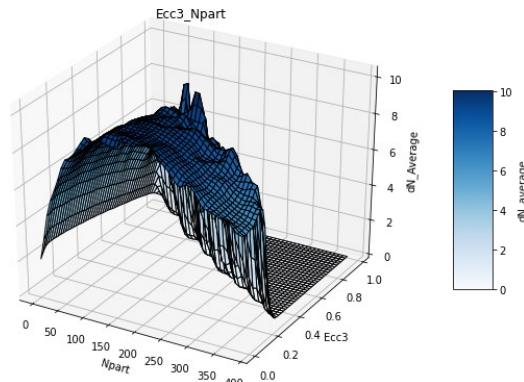


(b)  $\epsilon_3$  and  $N_{part}$

**Figure 7.** Contour plots and histograms illustrating the relationships of  $\epsilon_2$  with  $N_{part}$  and  $\epsilon_3$  with  $N_{part}$  for datasets where  $dN > 15$ .



(a) Three-dimensional depiction showcasing the relationship between  $\epsilon_2$  and  $N_{part}$ , with the mean  $dN$  functioning as the vertical axis.



(b) Three-dimensional visualization depicting the association between  $\epsilon_3$  and  $N_{part}$ , with the mean  $dN$  acting as the vertical axis.

**Figure 8.** Three-dimensional surface plots.



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