

Multiqubit quantum dense coding

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Abstract. Multipartite quantum-information transmission has been a key task for practice purposes and applications in the past decades. In this paper, we present a novel protocol for the multiqubit super dense coding, by using the Greenberger-Horne-Zeilinger (GHZ) states. Our protocol involves measurements doable in the current technique of labs, and thus is possibly realizable and may be extendable to more complicated cases for many systems. In quantum physics, the way the object is measured will affect the code, and the information won't work like puzzles as the condition in general physics which means that, when the divided message is put together, the complete information will be obtained. It requires more complex work to get useful information, like the cooperation of code accepters. To compare with the way in general physics, it is more safe in quantum physics.

Keywords: quantum information, multiqubit state, GHZ, dense coding.

1. Introduction

Quantum information processing has been a key task in applying quantum mechanics in recent decades. Ref. [1] introduced quantum communication using measurements on entangled Bell states. In Ref [2], the authors studied Teleporting an Unknown Quantum State through Dual Classical and Einstein-Podolsky-Rosen Channels. In Ref [3], what kinds of Fundamental limits will repeaterless quantum communications have been introduced by authors? The authors talked about a photonic integrated quantum secure communication system In Ref [4]. In Ref [5], the author discussed the method of quantum teleportation, about Ground-to-satellite. In Ref [6], light-to-motion Quantum teleportation was studied by the author. In Ref [7], N. Fiaschi and his team introduced Optomechanical quantum teleportation. In Ref [8], T. M. Graham's team used hyperentangled photons for Superdense teleportation. In Ref [9], the way of Probabilistic implementation operation by using a nonminimally entangled state introduced by L. Chen and Y.-X. Chen. In Ref [10], A. S. Cacciapuoti and his team studied quantum teleportation for the quantum internet, as entanglement meets classical communications. In Ref [11], a study of security during quantum dense coding in high-dimensions was introduced by Y.-X. Zhang's team. In Ref [12], A. Fonseca studied when noisy environments, and the condition of high-dimensional quantum teleportation. In Ref [13], a team of F. Shi introduced information masking in k-uniform quantum. In Ref [14], Y.-H. Luo's team studied in high dimensions, what quantum teleportation will be like. In Ref [15], a team of A. Barenco introduced quantum computation's Elementary gates.

2. Preliminaries

Here are some useful explanations for the mathematics that will be used. In this section, we introduce some useful explanations of the mathematics and notations that will be used in this paper. In Sec. 2.1, we review matrix basics and Kronecker products. Then we introduce the basic knowledge from quantum mechanics such as quantum states, Ket-bra notations, and entanglement. In Sec. 2.3, we introduce the fundamental idea of standard dense coding protocol for two systems.

2.1. Matrix basics

For matrix A $n \times p$, and matrix B $m \times q$,

$$A \otimes B = \begin{bmatrix} a_{1,1}B & \cdots & a_{1,p}B \\ \vdots & \ddots & \vdots \\ a_{n,1}B & \cdots & a_{n,p}B \end{bmatrix}. \quad (1)$$

Generally,

$$A \otimes B \neq B \otimes A, \quad (2)$$

Except when a is a scalar:

$$a \otimes A = A \otimes a = aA, \quad (3)$$

Or if a and b are vectors.

$$a^T \otimes b = ba^T = b \otimes a^T. \quad (4)$$

Furthermore, these formulas can be proved:

$$A \otimes (B \otimes C) = (A \otimes B) \otimes C. \quad (5)$$

$$A \otimes (B+C) = (A \otimes B) + (A \otimes C). \quad (6)$$

$$(A+B) \otimes C = (A \otimes C) + (B \otimes C). \quad (7)$$

For a matrix U , it has

$$U^\dagger U = U U^\dagger = I_n \quad (8)$$

2.2. quantum mechanics basics

When writing $|0\rangle$, which is called ket zero, it has: $= \begin{bmatrix} 1 \\ 0 \end{bmatrix}$.

$$|0\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}. \quad (9)$$

Similarly, for $|1\rangle$:

$$|1\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}. \quad (10)$$

If want to calculate $|0\rangle \otimes |0\rangle$, which is always be written as $|00\rangle$, according to (11):

$$|00\rangle = |0\rangle \otimes |0\rangle = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix}. \quad (11)$$

To calculate $|1\rangle \otimes |1\rangle$, similarly, written as $|11\rangle$, $= |1\rangle \otimes |1\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}$

$$|11\rangle = |1\rangle \otimes |1\rangle = \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix}. \quad (12)$$

For information $\frac{|00\rangle + |11\rangle}{\sqrt{2}}$, when it be polluted as $\cos \theta |00\rangle + \sin \theta |11\rangle$, $\theta \in [0, \frac{\pi}{2}]$, the expectation of getting the correct message from the polluted message is

$$E(\varphi(\theta)) = -(\cos \theta)^2 \log_2 (\cos \theta)^2 - (\sin \theta)^2 \log_2 (\sin \theta)^2. \quad (13)$$

Assume that there exists a space with length x , breadth y , and height z , within time t , and a particle is moving inside of it, the cumulation of probability will be 1.

2.3. dense coding of two parties

In this subsection, we introduce the knowledge and terminology used in this paper. We also review the standard protocol of dense coding. Source S generates an EPR pair shared by Alice and Bob, who can have any distance, for example, which is defined as:

$$|\phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle). \quad (14)$$

To get this,

$$CNOT(H \otimes I)|00\rangle = |\phi^+\rangle. \quad (15)$$

It can be shown as:

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & 1 \\ 1 & 0 & -1 & 0 \\ 0 & 1 & 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 0 \\ 0 \\ -1 \end{bmatrix}. \quad (16)$$

The graph of the Quantum Circuit is like this:

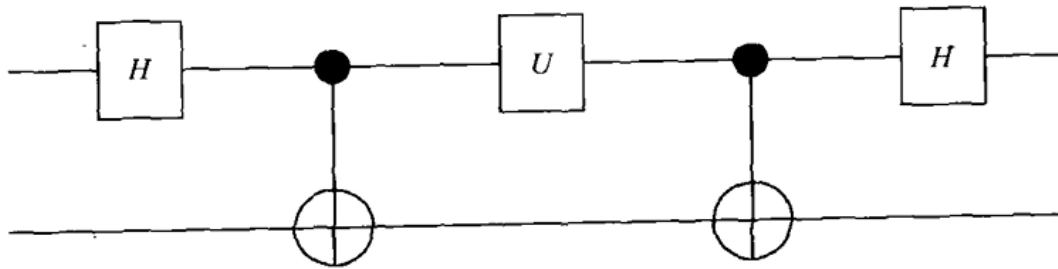


Figure 1. graph of the Quantum Circuit

Alice wants to send two bits to Bob, which could have 4 possible forms: 00, 11, 01, 10. To do this, she needs to separately do the following Unitary operations:

$$U = I, i\sigma_y, \sigma_x, \sigma_z \quad (17)$$

For example, if she wants to pass 01, the operation is:

$$\sigma_x \otimes I |\phi^+\rangle = |\psi^+\rangle. \quad (18)$$

In the matrix, it can be shown as:

$$\begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 0 \\ 1 \\ 1 \\ 0 \end{bmatrix}. \quad (19)$$

To learn about what Alice passed, Alice needs to pass her EPR pair to Bob at first, but if Robert, for example, gets only Bob's pair, without Alice's, when he tries to learn about the whole system, suppose $e_0 = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$, $e_1 = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$, $|a\rangle = |\phi^\pm\rangle, |\psi^\pm\rangle$; and $|a\rangle = (\langle a|)^\dagger$, where $|a\rangle$ is the entangled state shared by both the pairs. Robert will always obtain

$$Tr_A |a\rangle \langle a|_{AB} = \sum_{i=0}^1 (I_2 \otimes e_i^T) |a\rangle \langle a| (I_2 \otimes e_i) = \rho_B = \frac{1}{2} I_2. \quad (20)$$

By the way, if Alice's EPR pair was obtained by Robbert, for example, When he tries to decode it, he always obtains ρ_A equal to $\frac{1}{2} I_2$ in terms of (7).

So the formal operation is: Bob transforms the Bell state into the computational basis states, could be written as:

$$CNOT(H \otimes I)^{-1} = (H \otimes I)CNOT. \quad (21)$$

In matrix:

$$B = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & -1 & 0 \end{bmatrix}. \quad (22)$$

So it is easy to find that:

$$\begin{aligned} B|\psi^+\rangle &= |01\rangle, B|\psi^-\rangle = |11\rangle, \\ B|\phi^+\rangle &= |00\rangle, B|\phi^-\rangle = |10\rangle. \end{aligned} \quad (23)$$

In general physics, this dense coding is impossible because no matter how it is measured, the coded object will stay the same. Also, in this condition, when the divided message is put together, the complete information will be obtained, like puzzles. However, in quantum physics, the way the object is measured will affect the code, and the information won't work like puzzles.

3. Multiqubit Dense Coding

In this section, we introduce a generalized protocol of dense coding in which Alice sends messages to both Bob and Charlie.

Considering a situation if S generates a three EPR pair shared to Alice, Bob, and Charlie, who can have any distance, for example, which is defined as:

$$|\Psi^+\rangle = \frac{1}{\sqrt{2}} (|000\rangle + |111\rangle).$$

Depending on which kind of information Alice wants to pass, she can make following operations:

$$I_2 \otimes I_4 |\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|000\rangle + |111\rangle),$$

$$\sigma_x \otimes I_4 |\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|100\rangle + |011\rangle),$$

$$\sigma_z \otimes I_4 |\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|000\rangle - |111\rangle),$$

$$i\sigma_y \otimes I_4 |\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|011\rangle - |100\rangle),$$

After Alice making the similar operation like before, when Robert get Bob and Charlie's EPR pair, without Alice's. And when he tries to get the whole system, similar as formula 20, Robert will always get:

$$Tr_A |a\rangle\langle a|_{ABC} = \sum_{i=0}^1 (e_i^T \otimes I_4 \otimes I_4) |a\rangle\langle a| (e_i \otimes I_4 \otimes I_4) = \rho_{BC}.$$

The expression of ρ_{BC} depends on that of $|a\rangle$. If Robert takes $|a\rangle$ as $|a\rangle = \frac{1}{\sqrt{2}}(|000\rangle \pm |111\rangle)$, then $\rho_{BC} = \frac{1}{2} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$. If he takes $|a\rangle$ as $|a\rangle = \frac{1}{\sqrt{2}}(|100\rangle \pm |011\rangle)$, then $\rho_{BC} = \frac{1}{2} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{bmatrix}$. So he can't find out which one is the actually value of $|a\rangle$ according to the ρ_{BC} he get.

Forget about Robert. When Bob gets Alice's pair after Alice finishes her operation, which keeps the pair as $|\Psi^+\rangle$, for example, he doesn't need to get Charlie's pair, he just needs to communicate with Charlie to learn about the condition.

Bob can use the machine in the lab to measure the two pairs he gets. He may get $P_1 = |00\rangle\langle 00| + |11\rangle\langle 11|$, or $P_2 = |01\rangle\langle 01| + |10\rangle\langle 10|$, no matter what P_1 and P_2 is,

$$P_1 + P_2 = I_4.$$

We have two cases (i) and (ii), namely Bob obtains the measurement results P_1 and P_2 , respectively.

(i) Assume that Bob gets P_1 , and as he has no idea if the pair is equal to what, he will do the following operation, for example, he takes the pair equal to $\frac{1}{\sqrt{2}}(|100\rangle \pm |011\rangle)$:

$$(P_1 \otimes I_4) \frac{1}{\sqrt{2}}(|100\rangle \pm |011\rangle) = 0.$$

As the result is equal to 0, the pair will not be equal to $\frac{1}{\sqrt{2}}(|100\rangle \pm |011\rangle)$. But if he tries $\frac{1}{\sqrt{2}}(|000\rangle + |111\rangle)$, or $\frac{1}{\sqrt{2}}(|000\rangle - |111\rangle)$, he will get:

$$(P_1 \otimes I_4) \frac{1}{\sqrt{2}}(|000\rangle + |111\rangle) = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle),$$

$$(P_1 \otimes I_4) \frac{1}{\sqrt{2}}(|000\rangle - |111\rangle) = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \otimes \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle).$$

So, it can be one of $\frac{1}{\sqrt{2}}(|000\rangle + |111\rangle)$, or $\frac{1}{\sqrt{2}}(|000\rangle - |111\rangle)$. To learn about which one is the real value, Bob once again sent his pairs to the machine, then he will get one value from Q_1, Q_2, Q_3 . Similarly, no matter how,

$$Q_1 + Q_2 + Q_3 = I_4.$$

Assume that

$$Q_1 = \frac{|00\rangle\langle 00| + |11\rangle\langle 11|}{\sqrt{2}},$$

$$Q_2 = \frac{|00\rangle\langle 00| - |11\rangle\langle 11|}{\sqrt{2}},$$

$$Q_3 = |01\rangle\langle 01| + |10\rangle\langle 10|.$$

One can verify that Bob cannot obtain Q_3 due to zero probability. So we focus on two cases, Q_1 and Q_2 .

If Bob gets Q_1 , then he obtains

$$(Q_1 \otimes I_2)_{AB \otimes C} |a_I\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle + |1\rangle) = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$

$$(Q_1 \otimes I_2)_{AB \otimes C} |a_2\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle - |1\rangle) = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix},$$

After that he communicates with Charlie to learn about what the state is.

If Bob gets Q_2 , then he obtains

$$(Q_2 \otimes I_2)_{AB \otimes C} |a_I\rangle = \frac{|00\rangle - |11\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle - |1\rangle) = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix},$$

$$(Q_2 \otimes I_2)_{AB \otimes C} |a_2\rangle = \frac{|00\rangle - |11\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle + |1\rangle) = \frac{|00\rangle + |11\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$

After that he communicates with Charlie to learn about what the state is.

(ii) Assume that bob get $P_2 = |01\rangle\langle 01| + |10\rangle\langle 10|$, and as he has no idea of the pair is equal to what, he will do the following operation, for example, he takes the pair equal to $\frac{1}{\sqrt{2}}(|000\rangle \pm |111\rangle)$:

$$(P_2 \otimes I_4) \frac{1}{\sqrt{2}} (|000\rangle \pm |111\rangle) = 0.$$

As the result is equal to 0, the pair will not be equal to $\frac{1}{\sqrt{2}}(|000\rangle \pm |111\rangle)$. But if he tries $\frac{1}{\sqrt{2}}(|100\rangle + |011\rangle)$, or $\frac{1}{\sqrt{2}}(|100\rangle - |011\rangle)$, he will get

$$(P_2 \otimes I_4) \frac{1}{\sqrt{2}} (|100\rangle + |011\rangle) = \frac{1}{\sqrt{2}} (|10\rangle + |01\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle),$$

$$(P_2 \otimes I_4) \frac{1}{\sqrt{2}} (|100\rangle - |011\rangle) = \frac{1}{\sqrt{2}} (|10\rangle + |01\rangle) \otimes \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle).$$

So, it can be one of $\frac{1}{\sqrt{2}}(|100\rangle + |011\rangle)$, or $\frac{1}{\sqrt{2}}(|100\rangle - |011\rangle)$. To learn about which one is the real value, Bob once again sent his pairs to the machine, then he will get one value from R_1, R_2, R_3 . Similarly, no matter how,

$$R_1 + R_2 + R_3 = I_4.$$

Assume that

$$R_1 = \frac{|01\rangle + |10\rangle\langle 01| + \langle 10|}{\sqrt{2}},$$

$$R_2 = \frac{|01\rangle - |10\rangle\langle 01| - \langle 10|}{\sqrt{2}},$$

$$R_3 = |00\rangle\langle 00| + |11\rangle\langle 11|.$$

One can verify that Bob cannot obtain

R_3 due to the zero probability, so we focus on R_1, R_2 :

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$$(R_2 \otimes I_2)_{AB \otimes C} |a_2\rangle = \frac{|10\rangle + |01\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle - |1\rangle) = \frac{|10\rangle + |01\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix},$$

After that he communicates with Charlie to learn about what the state is.

If Bob gets Q_2 , then he obtains

$$(R_1 \otimes I_2)_{AB \otimes C} |a_2\rangle = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle - |1\rangle) = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix},$$

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After that he communicates with Charlie to learn about what the state is.

We stress that, if Alice does not trust Bob, then Alice may ask Charlie to stop the communication with Bob so that Bob has no way to figure out the code from Alice. In this sense, the security of three-party dense coding is again guaranteed by the basic rules of quantum mechanics.

In the following, we extend the above scheme to a multiqubit case. For this purpose, we introduce the multipartite GHZ state as follows.

Considering a situation if S generates a multiple EPR pair shared to Alice, Bob, Charlie and etc., who can have any distance, for example, which is defined as:

$$|\Psi^+\rangle = \frac{1}{\sqrt{2}} (|0 \dots 0\rangle + |1 \dots 1\rangle).$$

Depending on which kind of information Alice wants to pass, she can make following operations:

$$I_2 \otimes I_{2^n} |GHZ\rangle = \frac{1}{\sqrt{2}} (|00 \dots 0\rangle + |11 \dots 1\rangle),$$

$$\sigma_x \otimes I_{2^n} |GHZ\rangle = \frac{1}{\sqrt{2}} (|100 \dots 0\rangle + |011 \dots 1\rangle),$$

$$\sigma_z \otimes I_{2^n} |GHZ\rangle = \frac{1}{\sqrt{2}} (|00 \dots 0\rangle - |11 \dots 1\rangle),$$

$$i\sigma_y \otimes I_{2^n} |GHZ\rangle = \frac{1}{\sqrt{2}} (|011 \dots 1\rangle - |100 \dots 0\rangle),$$

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After Alice making the similar operation like before, when Robert get everyone's EPR pair, but without Alice's. And when he tries to get the whole system, similar as formula 20, Robert can't find out which one is the actually value of $|a\rangle$ according to the $\rho_{BC\dots}$ he gets.

Forget about Robert. Set Bob's pair as B_1 , others as but not including Alice, as B_n . When Bob gets Alice's pair after Alice finishes her operation, which keeps the pair as $|\Psi^+\rangle$, for example, he doesn't need to get others pair, he just needs to communicate with them to learn about the condition.

Bob can use the machine in the lab to measure the two pairs he gets. He may get $P_1 = |000 \dots 0\rangle + |11 \dots 1\rangle$, or $P_2 = |01 \dots 0\rangle + |10 \dots 1\rangle$, no matter what P_1 and P_2 is,

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If Bob gets Q_2 , then he obtains

$$(R_2 \otimes I_2)_{AB \otimes C} |a_2\rangle = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle - |1\rangle) = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ -1 \end{pmatrix},$$

$$(R_2 \otimes I_2)_{AB \otimes C} |a_1\rangle = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} (|0\rangle + |1\rangle) = \frac{|10\rangle - |01\rangle}{\sqrt{2}} \frac{1}{2} \begin{pmatrix} 1 \\ 1 \end{pmatrix},$$

After that he communicates with others to learn about what the state is.

We stress that, if Alice does not trust Bob, then Alice may ask Charlie to stop the communication with Bob so that Bob has no way to figure out the code from Alice. In this sense, the security of three-party dense coding is again guaranteed by the basic rules of quantum mechanics.

4. Conclusion

We have presented a new protocol of multiqubit super dense coding, by using the multiqubit Greenberger-Horne-Zeilinger states. We have given details for the three-qubit case, and we extended it to the multiqubit case in a more succinct way. The next step in this paper is to study whether one can similarly construct more quantum-information processing tasks by using a similar idea, such as multipartite quantum secret sharing and masking.

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